



Including factorisation and renormalisation scale -choice and -variation into FastNLO



The FastNLO idea

*fast*NLO

- FastNLO factorizes the cross section calculation for an a-posteriori inclusion of pdf's and alpha_s for jet-production
- The basic cross section formula for pQCD (DIS)

$$\sigma = \sum_{a,n} \int_0^1 dx \alpha_s^n(\mu_r) \cdot c_{a,n}\left(\frac{x_{Bj}}{x}, \mu_r, \mu_f\right) \cdot f_a(x, \mu_f)$$

- n: order alpha_s (perturbative order)
- a: number of flavors (internally it is Δ, g, Σ)
- c: perturbative coefficients from the matrix elements
- f: pdf



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- the 'x'-integral is getting factorized

$$f_a(x) \cong \sum_i f_a(x_i) \cdot E^{(i)}(x)$$

- and calculated at some 'eigen'values

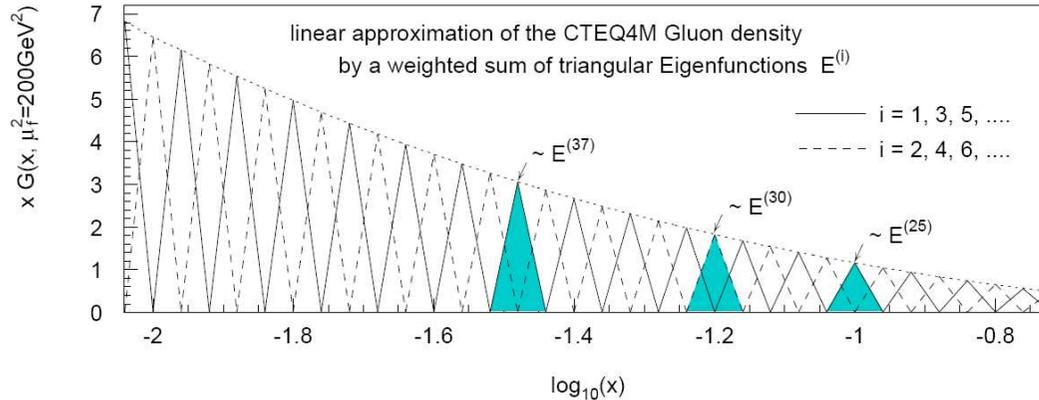
$$\sigma \cong \sum_{a,n,i} \alpha_s^n f_a(x_i) \underbrace{\int dx c_{a,n}\left(\frac{x_{Bj}}{x}\right) E^{(i)}(x)}_{\tilde{\sigma}} = \sum_{a,n,i} \alpha_s^n f_a(x_i) \tilde{\sigma}_{a,n}^{(i)}$$

BUT:
what about
the scales ?



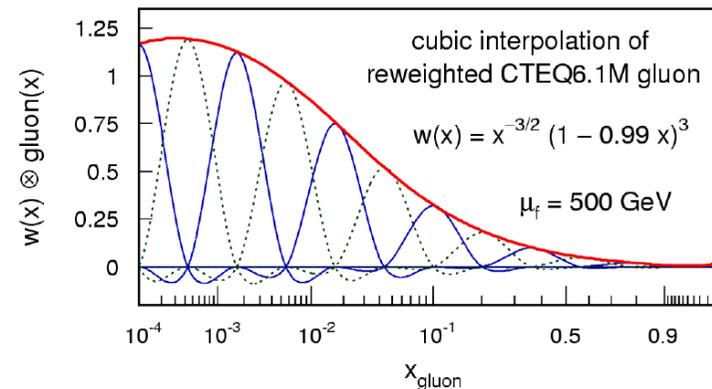
approximation of 'gluon' density

- as example: approximation by triangular eigenfunctions



$$E^{(i)}(x) = \begin{cases} 1 & x = x_i \\ \log(x_{i-1}) = \log(x) & x_{i-1} < x < x_i \\ \log(x_{i+1}) = \log(x) & x_i < x < x_{i+1} \\ 0 & x < x_{i-1}, \text{ or } x > x_{i+1} \end{cases}$$

- FastNLO uses more sophisticated cubic eigenfunctions



- Basically: Each 'weight' is distributed over 4 scale nodes
- Afterwards, we only have to sum up those nodes



The scale dependence

- α_s and pdf have to be evaluated at the correct μ value
- FastNLO takes $\mu_r = \mu_f = \mu$
- FastNLO calculates cross section within 'sufficiently' small intervals of μ
 - those intervals are also interpolated using cubic eigenfunctions

$$\sigma = \sum_a^{\text{flavors}} \sum_n^{\text{orders}} \sum_i^{x\text{-nodes}} \sum_j^{\mu\text{-nodes}} \alpha_s^n(\mu(j)) \cdot \tilde{\sigma}_{a,n,j}^{(i)}(\mu(j)) \cdot f_a(x_i, \mu(j))$$

- Save another 'table' of μ -nodes $\mu(j)$
- Those 'tables' are calculated for each bin of our measurement



Needs for improvements

- $\mu_r = \mu_f$ might not be wanted (e.g. low Q^2 jets)
- Some studies for different scales might be done e.g.
 - Q^2
 - P_T jet in Breit frame
 - $\langle p_T \rangle$ all jets
 - $\sqrt{(Q^2 + p_T^2)/2}$
 - $\sqrt{(Q^2 + p_T^2)/4}$
 - $\sqrt{(Q^2 + p_T^2)}$
- Scales should be consistent for combined fits
 - Zeus, H1, should not choose different scales in pdf fits
- Estimation of theory error
 - scales are varied by factor 2 up and down
 - BUT: $\tilde{\sigma}$ is scale dependent
 - Scale variations are only approximations
- Solution:

We store scale independent weights, and do all the scale calculations when calculating the cross section



cross section in nlojet++

- The cross section calculation in some detail

$$\sigma = \sigma^{born} + \sigma^{real} + \sigma^{virtual}(\mu_r, \mu_f)$$

$$\sigma^{virtual} = \sigma^{sub} + \sigma^{finix}(\mu_r, \mu_f) + \sigma^{fini-1}(\mu_r, \mu_f)$$

- while the cross section is only a sum of weights in the MC integration

$$w = W_{PS} \cdot pdf \cdot \alpha_s \cdot M(\mu_r, \mu_f) \cdot \alpha_{em}^2(Q^2) \cdot (\hbar c)^2$$

$$M^{finite} = M_0 + M_f \cdot \ln\left(\frac{\mu_f^2}{Q^2}\right) + M_r \cdot \ln\left(\frac{\mu_r^2}{Q^2}\right)$$

- which is adapted in fastNLO like

$$w = W_{PS} \cdot \frac{1}{x} \cdot 1 \cdot M(\mu_r, \mu_f) \cdot \alpha_{em}^2(Q^2) \cdot (\hbar c)^2 \equiv c \cdot M$$



improved FastNLO

- Access M_0 , M_f , and M_r directly
 - slight hack of nlojet++
- Store THREE tables for
 - $M_0 * c \rightarrow \sigma()$
 - $M_f * c \rightarrow \sigma_f(\mu_f)$
 - $M_r * c \rightarrow \sigma_r(\mu_r)$
 - do the multiplication the same way as nlojet++ does it

- Calculating the cross section

$$\sigma = \sigma_0 + \sigma_f \cdot \ln\left(\frac{\mu_f^2}{Q^2}\right) + \sigma_r \cdot \ln\left(\frac{\mu_r^2}{Q^2}\right)$$

- remember: σ_{xy} here is actually a sum over $x, n, a, (i)$
- Need for knowledge of Q^2 , μ_r and μ_f
 - -> Additional "Scale" table for Q^2
- Better
 - Additional "Scale" table for Q AND one table for p_T !!!



new final formula

- FastNLO formula with three scale independent tables σ_0 , σ_f , σ_r and two 'scale' lookup tables (q,p)

$$\sigma_{Bin} = \sum_a^{\text{flavors}} \sum_n^{\text{orders}} \sum_i^{x\text{-nodes}} \sum_q^{Q\text{-nodes}} \sum_p^{p_T\text{-nodes}} \left(\begin{aligned} & \alpha_s^n(\mu_r(q,p)) \cdot \sigma_{0a,n}^{(i)} \cdot f_a(x_i, \mu_f(q,p)) + \\ & \alpha_s^n(\mu_r(q,p)) \cdot \sigma_{fa,n}^{(i)} \cdot \ln\left(\frac{\mu_f(q,p)^2}{Q(q)^2}\right) \cdot f_a(x_i, \mu_f(q,p)) + \\ & \alpha_s^n(\mu_r(q,p)) \cdot \sigma_{ra,n}^{(i)} \cdot \ln\left(\frac{\mu_r(q,p)^2}{Q(q)^2}\right) \cdot f_a(x_i, \mu_f(q,p)) \end{aligned} \right)$$

- while μ_r and μ_f are just some functions of the stored (in the 'scale table') Q and p_T values at p and q
 - e.g. $\mu_r(q,p) = \text{sqrt}((Q_q^2 + p_{T,p}^2)/2)$
- one could also store sth. different than p_T but one always needs Q^2



pp and ppbar

- somehow simpler and more complicated
 - sum over x_1 and x_2 for both hadrons
 - 7(6) types of pdf – linear combinations
 - 6 types of contributions to xs formula
 - born, real, sub, finix1, finix2, fini1
- μ_r and μ_f dependence is more trivial

$$\sigma = \sigma_0 + \sigma_f \cdot \log(\mu_f^2) + \sigma_r \cdot \log(\mu_r^2)$$

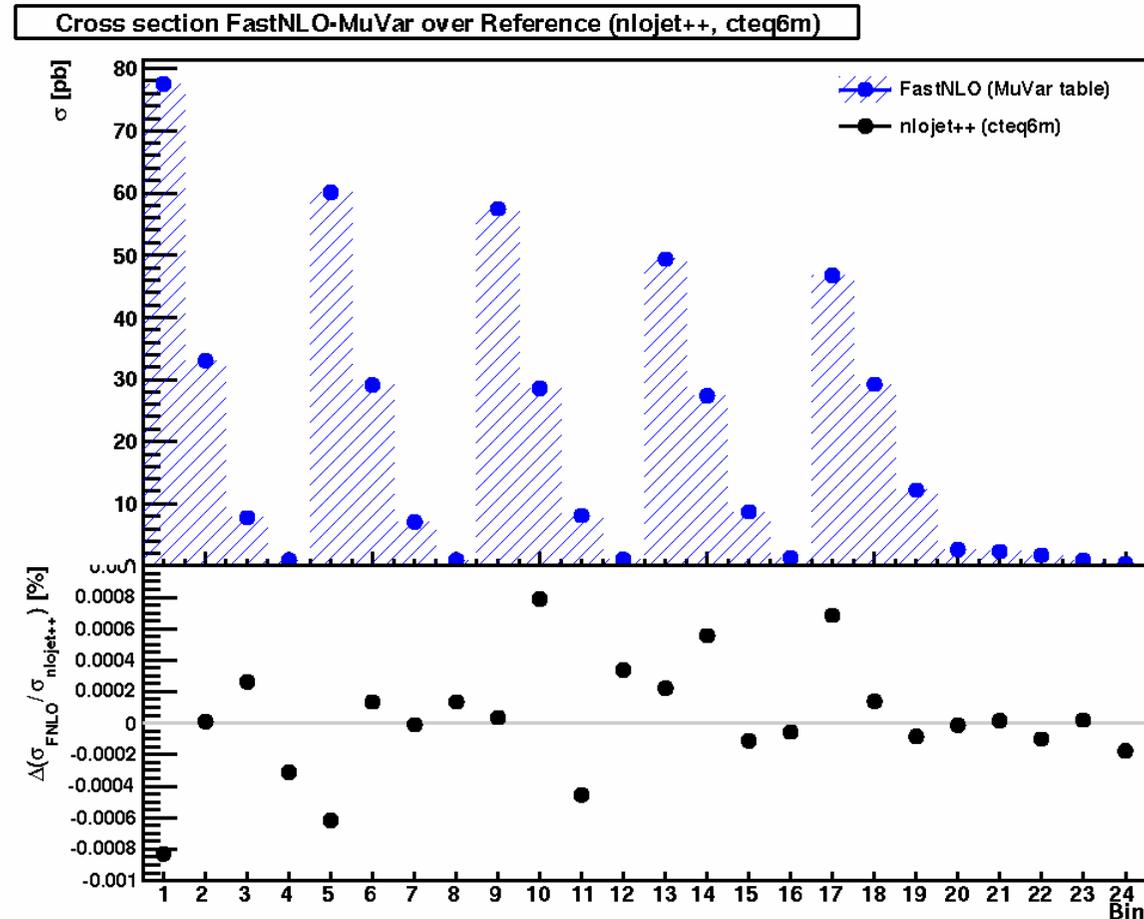
- final formula is rather simple then



Make a long story short



- Testing with inclusive jets at high Q^2
- choosing
 - 70 x-nodes
 - 30 nodes for Q
 - 30 nodes for p_T
- Using always cubic interpolation
 - each weight is distributed over $4 \times 4 \times 4$ nodes
- choosing (!!) $\mu_r = \mu_f = Q^2$
- Tablesize: 200 MB
- **FNLO precision**
 - **better than: 0.0005%**

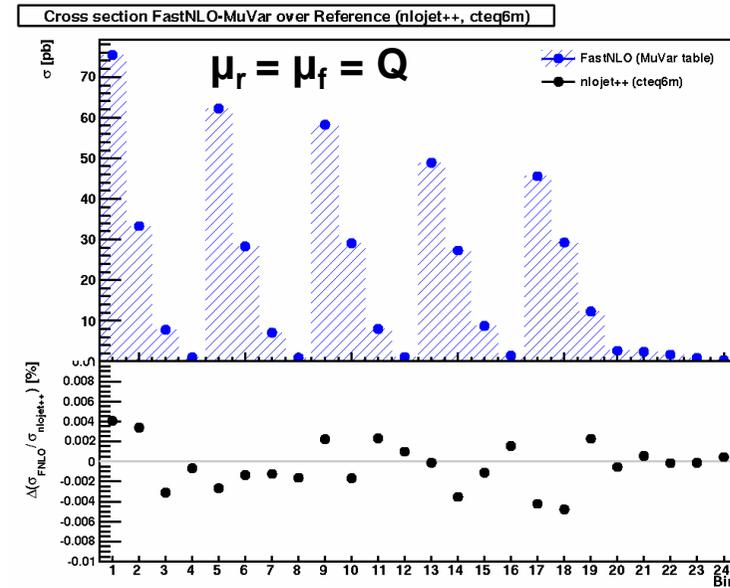




Use less bins, and study scale choice



- choosing
 - 30 x-nodes
 - 20 nodes for Q
 - 15 nodes for pT
 - size: 30MB
- use the appropriate reference cross section
- Precision is mostly dependent on number of x-bins:
 - "High Q², High pT" bins span over smaller x-range -> higher precision with fixed number of bins (Backup)
 - Different number of x-nodes per bin is already implemented!

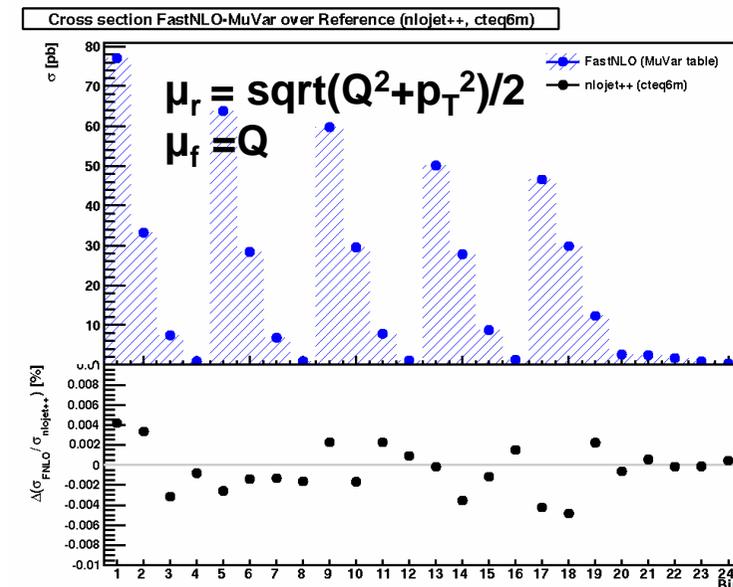
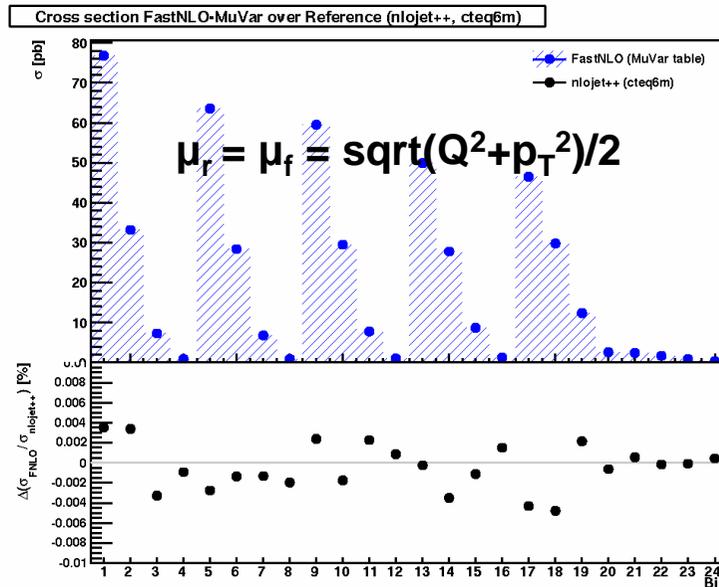
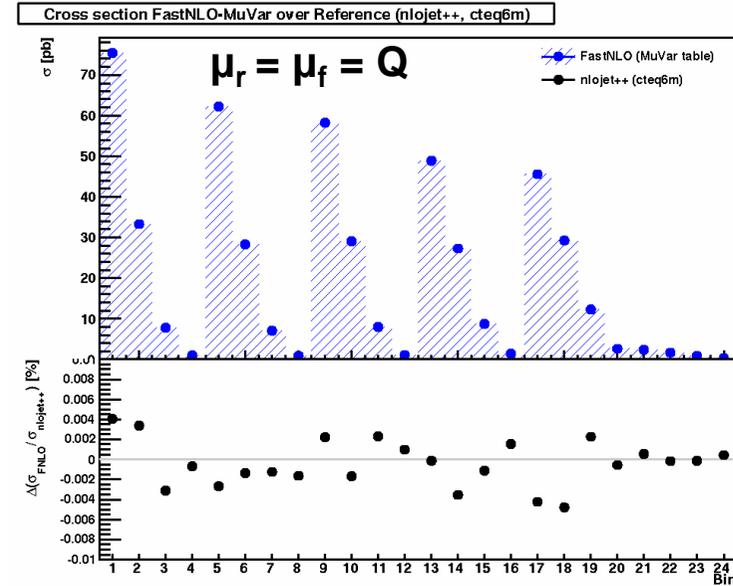




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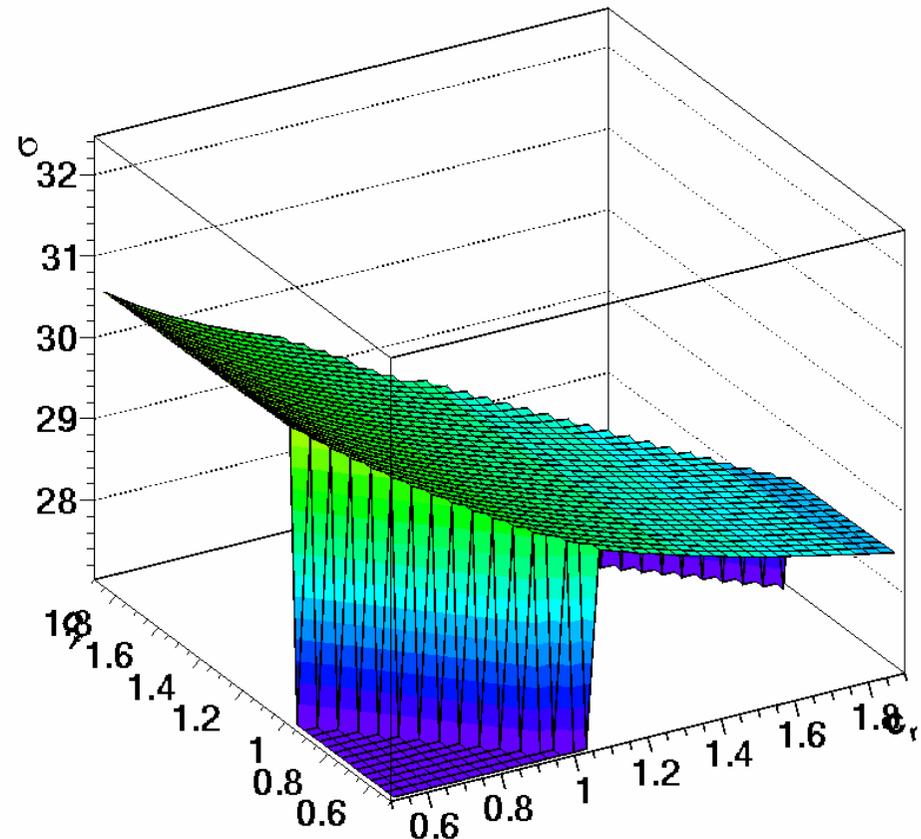


Scale variation for 'theory error'



- functional form of scales is free to choose
 - also $c_r \cdot \mu_r$ and $c_f \cdot \mu_f$ is a valid choice
- studying scale variation with very high precision and very high statistics!
 - constrain:
 $0.5 < c_r \cdot c_f < 2.0$
- new and very precise determination of conventionally defined 'theory error'

Cross section bin 9





Comparision



	FastNLO (to be v2.0)	This FastNLO mod
Size of table (speed)	$n_{\text{Bins}} \times n_x \times n_{\text{mur}} [\times n_{\text{muf}}] \times n_{\text{scalevar}} \times n_{\text{proc}} \times n_{\text{ord}}$	$n_{\text{Bins}} \times n_x \times n_{\text{Pt}} \times n_{\text{Q}^2} \times n_{\text{proc}} \times n_{\text{ord}} \times 3$
choice of renormalization scale	Fixed when creating table e.g. $\mu_r^2 = (Q^2 + p_T^2)/2$	Any function using Q^2 and P_t is possible (or any other variables used, when creating table)
choice of factorization scale	equals renormalization scale (v2.0) (I have a mod, where $\mu_r \neq \mu_f$)	Any function using Q^2 and P_t is possible (or any other variables used, when creating table)
variation of μ_r (e.g. $\times 2$, $\times 0.5$)	fixed, when creating table	every factor is possible
variation of μ_f (e.g. $\times 2$, $\times 0.5$)	fixed and must be equal μ_r	every factor is possible (independent of μ_r)
free variation of scale factor	μ_r : only in an approximation μ_f : not supported	see above
How to determine theory uncertainty	vary μ_r and μ_f up and down simultaneously by $\times 2$ and $\times 0.5$	vary μ_r and μ_f independently with $0.5 < c_r^* c_f < 2.0$ and look for highest and lowest prediction (e.g. ATLAS jets: arXiv:1009.5908)



Summary

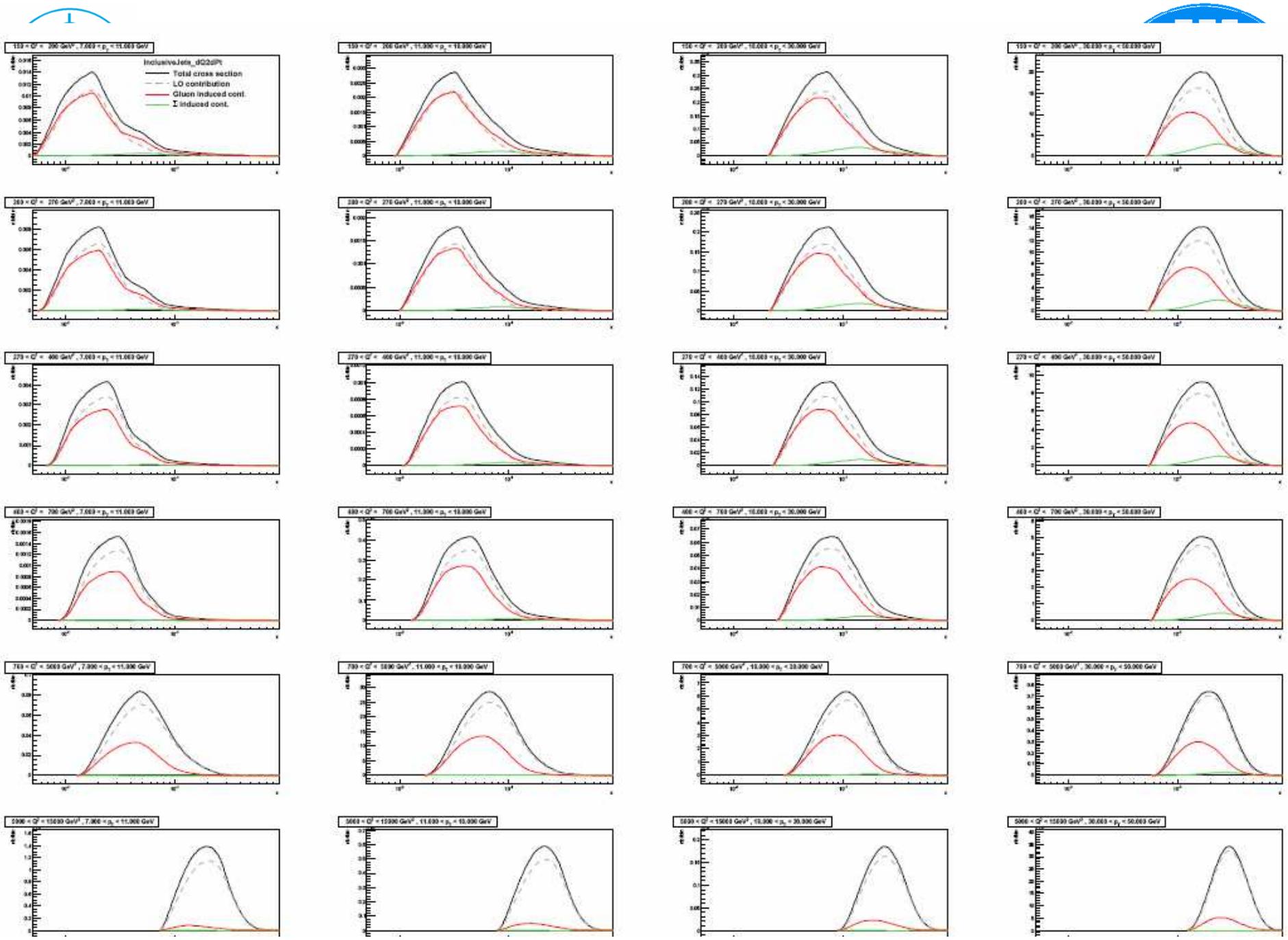


- The FastNLO concept was improved
 - Scale independent tables allow to choose renormalization and factorization scale afterwards almost arbitrary
 - a third variable (like Θ , pt-leading) could be implemented
 - Scale variations for theory error determination can be performed much more sophisticated and without any approximation
 - The precision of this FastNLO version can reach more than 0.0005%
 - This method is implemented and well tested for DIS processes
 - It is already implemented for pp and ppbar but testing is needed (reader is missing)
-
- Authors of FastNLO did not answer my last eMail about that
 - I would propose to use this method for our upcoming α_s fits
 - Method is also great for combined fits and cross checks with ZEUS, if they use different scales than we do



Backup





Daniel Britzger – FastNLO upgrade